# Seminvariants for Non-centrosymmetric Space Groups with Conventional Centered Cells 

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#### Abstract

The relationship of phase to the choice of origin, enantiomorph or frame of reference is clarified for those non-centrosymmetric space groups for which the conventional unit cell is not primitive. The theory employs special lincar eombinations of the phases, the structure seminvariants. Simple procedures are developed for selecting the origin by first fixing the functional form of the structure factor, then specifying the sign of a seminvariant when required, and, finally, specifying arbitrarily the values of a suitable set of phases.

This paper completes the study of the seminvariants for all the space groups.


## 1. Introduction

In the direct determination of phases from the obverved intensities it is necessary to relate the values of the phases to the choice of origin, reference frame and enantiomorph. This problem has already been treated for the centrosymmetric space groups (Hauptman \& Karle, 1953, 1959) and for the non-centrosymmetric space groups for which the conventional unit cell is primitive (Hauptman \& Karle, 1956). It was found that certain linear combinations of the phases, the structure seminvariants, play a fundamental role in these studies. The seminvariants show which linear combinations are determined by the intensities alone and how specifications of phases are to be made to fix the origin, frame and enantiomorph.

In this paper we complete the study of seminvariants for the various space groups, by considering the noncentrosymmetric space groups for which the conventional unit cell is non-primitive. The non-primitive cell is transformed to an appropriate primitive cell by means of well-known transformations. The methods referred to above are then immediately applicable.

## 2. Primitive unit cells

The coordinates representing the space group relative to a primitive unit cell are obtained from those corresponding to a non-primitive unit cell (International Tables, 1952) by means of the following matrices:

$$
\begin{align*}
& C \rightarrow P,\left(\begin{array}{rrr}
1 & 1 & 0 \\
1 & -1 & 0 \\
0 & 0 & -1
\end{array}\right), \\
& A \rightarrow P,\left(\begin{array}{rrr}
-1 & 0 & 0 \\
0 & -1 & 1 \\
0 & 1 & 1
\end{array}\right), \\
& I \rightarrow P,\left(\begin{array}{rrr}
0 & 1 & 1 \\
1 & 0 & 1 \\
1 & 1 & 0
\end{array}\right),
\end{align*}
$$

$$
F \rightarrow P,\left(\begin{array}{rrr}
-1 & 1 & 1 \\
1 & -1 & 1 \\
1 & 1 & -1
\end{array}\right)
$$

The results are shown in Table 1.

## 3. Definitions

In the discussion to follow several concepts will be employed, namely, linear and rational dependence and independence, primitive sets, equivalence and seminvariance. These concepts are defined and developed in our previous papers (Hauptman \& Karle, 1956, $\$ 3-\S 7 ; 1959, \S 4$ and $\$ 5)$ to which the reader is referred. They culminate in the main result which identifies the structure seminvariants, namely those linear combinations

$$
\sum_{\mathrm{h}} A_{\mathrm{h}} \psi_{\mathrm{h}},
$$

where the $A_{h}$ are integers satisfying

$$
\sum_{\mathbf{h}} A_{\mathbf{h}} \mathbf{h}_{s} \equiv 0\left(\bmod \omega_{s}\right)
$$

$\mathbf{h}_{s}$ is the vector seminvariantly associated with the phase $\psi_{h}$, and $\omega_{s}$ is the seminvariant modulus of the type. The seminvariant vectors and moduli are readily derived from the equivalence classes. These are listed in Table 2. It should be noted that the functional form of the structure factor is the same for all origins comprising an equivalence class.

This paper is concerned with describing in detail simple methods for selecting the origin in each of fourteen types of space groups. The procedures to be presented are of a relatively simple nature, although more general procedures may be readily derived from Table 2.

All the theorems of this paper are valid under either one of the following hypotheses.

Hypothesis A: The crystal structure is given; or
Hypothesis $B$ : A sufficiently large number of struc-

Table 1. Coordinates for centered non-centrosymmetric space groups referred to a primitive unit cell


Table 1 (cont.)

ture-factor magnitudes is given (so that, for a fixed functional form of the structure factor, the magnitudes of all the structure seminvariants are determined) and the sign of any one structure seminvariant, the magnitude of which is different from 0 or $\pi$, has been arbitrarily specified.

It is further assumed throughout this paper that the functional form of the structure factor is fixed.

## 4. The remaining types of space groups

## 4-01. Type $2 \mathrm{P}_{0} 02$

Theorem $4 \cdot 01 \cdot 1$. A single phase $\psi_{h}$ is a structure seminvariant, i.e. its value is uniquely determined if, and only if, $h=k$ and $l$ is even.

Theorem $4 \cdot 01 \cdot \underline{2}$. Let $h_{1}=k_{1}$. Then any phase $\psi_{h_{1}}$ which is linearly semi-independent (i.e. $l_{1}$ is odd) has
Table 2. Equivalence classes, seminvariant vectors and seminvariant moduli for the centered non-centrosymmetric space groups, referred to a primitive unit cell

just two possible values, and these differ from each other by $\pi$. Either one of these two values may be chosen. Once this is done, then the value of any phase $\varphi_{h}$ which is linearly semi-dependent on $\varphi_{h_{1}}$ (i.e. $h=k$ ) is uniquely determined.

Theorem $4 \cdot 01 \cdot 3$. Let $l_{2}$ be even. Then the value of any phase $\varphi_{\mathbf{h}_{2}}$ which is linearly semi-independent (i.e. $h_{2}-k_{2} \neq 0$ ) may be specified arbitrarily. Once this is done, the value of any phase $\mathscr{F}_{h}$ which is linearly semi-dependent on $\varphi_{\mathbf{h}_{2}}$ (i.e. $l$ is even and $h-k$ is divisible by $h_{2}-k_{2}$ ) is uniquely determined. Any phase $\psi_{\mathrm{h}}$ which is rationally semi-dependent on ${q_{\mathbf{h}_{2}}}$ is also linearly semi-dependent on $\varphi_{h_{2}}$, whence its value is uniquely determined, provided that $\varphi_{h_{2}}$ is semiprimitive, i.e. provided that $h_{2}-k_{2}= \pm 1$.

Theorem $4 \cdot 01 \cdot 4$. Let $h_{1}=k_{1}$ and $l_{2}$ be even. Let $\psi_{h_{1}}$ and $\varphi_{h_{2}}$ be any two phases which constitute a linearly semi-independent set (i.e. $l_{1}$ is odd and $h_{2}-k_{2}=0$ ). In accordance with the two previous theorems, either one of the two possible values of $\psi_{h_{1}}$ may be chosen while the value of $\varphi_{h_{2}}$ may be specified arbitrarily. Once this is done the value of any phase $q_{h}$ which is linearly semi-dependent on the pair $\varphi_{\mathbf{h}_{1}}, \varphi_{\mathbf{h}_{2}}$ is uniquely determined. Any phase $\phi_{h}$, of necessity rationally semi-dependent on the pair $\psi_{\mathbf{h}_{1}}, q_{\mathbf{h}_{2}}$, is also linearly semi-dependent on this pair, whence its value is uniquely determined, provided that the pair $\psi_{\mathbf{h}_{1}}, \psi_{\mathbf{h}_{2}}$ is semi-primitive, i.e. provided that $h_{2}-k_{2}= \pm 1$.

### 4.02. Type $2 P(0)$

Theorem $4 \cdot 0 \cdot 2 \cdot 1$. A single phase $\psi_{h}$ is a structure seminvariant, i.e. its value is uniquely determined, if, and only if, $h+k=l=0$.

Theorem $4 \cdot\left(02 \cdot 2\right.$. The value of any phase $\psi_{h_{1}}$ which is linearly semi-independent (i.e. $h_{1}+k_{1}$ and $l_{1}$ are not both zero) may be specified arbitrarily. Once this is done, the value of any phase $\varphi_{h}$ which is linearly semi-dependent on $\varphi_{\mathbf{h}_{1}}$ is uniquely determined. Any phase $\psi_{\mathbf{h}}$ which is rationally semi-dependent on $\psi_{h_{1}}$ is also linearly semi-dependent on $\varphi_{\mathbf{h}_{1}}$, whence its value is uniquely determined, provided that $\varphi_{\mathbf{h}_{1}}$ is semiprimitive, i.e. provided that the greatest common divisor of $h_{1}+k_{1}$ and $l_{1}$ is unity.

Theorem $4 \cdot 0 \cdot 2 \cdot 3$. The values of any two phases $\psi_{\mathbf{h}_{1}}, \psi_{\mathbf{h}_{2}}$, constituting a linearly semi-independent set, i.e.

$$
\left|\begin{array}{ll}
h_{1}+k_{1} & l_{1} \\
h_{2}+k_{2} & l_{2}
\end{array}\right| \neq 0
$$

may be specified arbitrarily. Once this is done, the value of any phase $\psi_{h}$ which is linearly semi-dependent on the pair $\varphi_{h_{1}}, \mathscr{F}_{h_{2}}$ is uniquely determined. Any phase $\psi_{h}$, of necessity rationally semi-dependent on the pair $q_{h_{1}}, \psi_{h_{2}}$, is also linearly semi-dependent on this pair, whence its value is uniquely determined, provided that the pair $\mathscr{F}_{\mathbf{h}_{1}}, \mathscr{F}_{\mathbf{h}_{2}}$ is semi-primitive, i.e. provided that

$$
\left|\begin{array}{ll}
h_{1}+k_{1} & l_{1} \\
h_{2}+k_{2} & l_{2}
\end{array}\right|= \pm 1
$$

AC 14-1:
4.(13. Type $2 P 20$

### 4.04. Type $2 P 22$

These types have been treated previously (Hauptman \& Karle, 1956).

### 4.05. Type $2 P_{1} 20$

Theorem $4 \cdot 0 \tilde{r} \cdot 1$. A single phase $\psi_{h}$ is a seminvariant if, and only if, $h$ is cven and $k+l=0$.

Theorem $4 \cdot 05 \cdot 2$. Let $k_{1}+l_{1}=0$. Then any phase $\psi_{h_{1}}$ which is linearly semi-independent (i.e. ${h_{1}}_{1}$ is odd) has just two possible values and these differ from each other by $\pi$. Either one of these two values may be chosen. Once this is done, the value of any phase $\psi_{h}$ which is linearly semi-dependent on $\mathscr{F}_{\mathbf{h}_{1}}$ (i.e. $k+l=0$ ) is uniquely determined.

Theorem $4 \cdot 05 \cdot 3$. Let $h_{2}$ be even. Then the value of any phase $\psi_{h_{2}}$ which is linearly semi-independent (i.e. $k_{2}+l_{2}=0$ ) may be specified arbitrarily. Once this is done, the value of any phase $q_{h}$ which is linearly semidependent on $\varphi_{\mathbf{h}_{2}}$ is uniquely determined. Any phase $\varphi_{h}$ which is rationally semi-dependent on $\varphi_{h_{2}}$ is also linearly semi-dependent on $\varphi_{h_{2}}$, whence its value is uniquely determined provided that $q_{h_{z}}$ is semi-primitive, i.e. provided that $k_{2}+l_{2}= \pm 1$.

Theorem $4 \cdot 0 \cdot 5 \cdot 4$. Let $k_{1}+l_{1}=0$ and $h_{2}$ be even. Let $\psi_{h_{1}}$ and $\psi_{h_{2}}$ be any two phases which constitute a linearly semi-independent set (i.e. $h_{1}$ is odd and $k_{2}+l_{2} \neq 0$ ). In accordance with the two previous theorems either one of the two possible values of $\psi_{h_{1}}$ may be chosen while the value of $\varphi_{\mathbf{h}_{2}}$ may be specified arbitrarily. Once this is done, the value of any phase $\psi_{\mathbf{h}}$ which is linearly semi-dependent on the pair $\psi_{\mathbf{h}_{1}}, \psi_{\mathbf{h}_{2}}$ is uniquely determined. Any phase $q_{\mathbf{h}}$, of necessity rationally semi-dependent on the pair $\psi_{\mathbf{h}_{3}}, \psi_{\mathbf{h}_{2}}$, is also linearly semi-dependent on this pair, whence its value is uniquely determined, provided that the pair $\psi_{\mathbf{h}_{1}}, \varphi_{\mathbf{h}_{2}}$ is semi-primitive, i.e. provided that $k_{2}+l_{2}= \pm 1$.

## 4-06. Type $2 P_{1} 022$

Theorem $4 \cdot 06 \cdot 1$. A single phase $\psi_{h}$ is a seminvariant if, and only if, $h+k=0$ and $h \equiv l(\bmod 2)$.

Theorem $4 \cdot 06 \cdot 2$. Let $h_{1}+k_{1}=0$. Then any phase $\psi_{\mathbf{h}_{1}}$ which is linearly semi-independent (i.e. $h_{1}+l_{1}$ is odd) has just two possible values and these differ from each other by $\pi$. Either one of these two values may be chosen. Once this is done, the value of any phase $\psi_{h}$ which is linearly semi-dependent on $\psi_{\mathbf{h}_{1}}$ is uniquely determined.

Theorem 4-06•3. Let $h_{2}+k_{2} \neq 0$, so that $\psi_{h_{2}}$ is linearly semi-independent. Then the value of $\phi_{\mathbf{h}_{2}}$ may be specified arbitrarily. Once this is done, the value of any phase $\psi_{h}$, which is linearly semi-dependent on $\psi_{h_{2}}$, is uniquely determined. Any phase $\psi_{h}$ which is rationally semi-dependent on $\mathscr{F}_{h_{2}}$ is also linearly semidependent on $\mathscr{f}_{h_{2}}$ provided that $q_{h_{2}}$ is semi-primitive, i.e. provided that $h_{2}+k_{2}= \pm 1$.

Theorem 4•06•4. Choose $q_{h_{1}}$ and $q_{h_{2}}$ as in the previous
two theorems. In accordance with these theorems either of the two possible values of $q_{\mathbf{h}_{1}}$ may be chosen and the value of $\varphi_{h_{2}}$ may be specified arbitrarily. Once this is done any phase $\psi_{\mathbf{h}}$, of necessity rationally semi-dependent on the pair $\varphi_{\mathbf{h}_{1}}, \varphi_{\mathbf{h}_{2}}$, is also linearly semi-dependent on this pair, whence its value is uniquely determined, provided that the pair $\varphi_{\mathbf{h}_{1}}, \varphi_{\mathbf{h}_{2}}$ is semi-primitive, i.e. provided that $h_{2}+k_{2}= \pm 1$.

### 4.07. Type $2 P_{1} 222$

Theorem $4 \cdot 07 \cdot 1$. A single phase $\varphi_{\mathbf{h}}$ is a seminvariant if, and only if, $h \equiv k \equiv l(\bmod 2)$.

Theorem $4 \cdot 07 \cdot 2$. Any phase $\psi_{\mathbf{h}_{1}}$ which is linearly semi-independent has just two possible values and these differ from each other by $\pi$. Either one of these two values may be chosen. Once this is done the value of any phase which is linearly semi-dependent on $\varphi_{h_{1}}$ is uniquely determined.

Theorem 4.07.3. Let the pair of phases $\varphi_{\mathbf{h}_{1}}, \varphi_{\mathbf{h}_{1}}$, be a linearly semi-independent set. In accordance with the previous theorem, either of the two possible values of $\varphi_{h_{1}}$ may be chosen and either of the two possible values of $\varphi_{\mathbf{h}_{2}}$ may be chosen. Once this is done, the value of any phase $\varphi_{h}$, of necessity linearly semidependent on the pair $\psi_{h_{1}}, \varphi_{h_{2}}$, is uniquely determined.

### 4.08. Type $3 \mathrm{P}_{2}$ 2

This type has been previously described (Hauptman \& Karle, 1956).

### 4.09. Type $3 P_{2} 4$

Theorem $4 \cdot 09 \cdot 1$. A single phase $\psi_{h}$ is a seminvariant if, and only if, $h+k+l \equiv 0(\bmod 4)$.

Theorem $4 \cdot 09 \cdot 2$. Let the phase $q_{h_{1}}$ be linearly semiindependent. Depending upon whether $h_{1}+k_{1}+l_{1}$ is odd or even, there are four or two possible values for $\tau_{h_{1}}$ (differing by $\pi / 2$ or $\pi$ respectively).

In the first case any of the four possible values for $7 \mathbf{h}_{1}$ may be chosen. Once this is done the value of any phase $\psi_{\mathbf{h}}$, of necessity linearly semi-dependent on $\gamma_{h_{1}}$, is uniquely determined.

In the second case either of the two possible values for $\varphi_{h_{1}}$ may be chosen. Once this is done then the value of any phase $\psi_{h}$ which is linearly semi-dependent on $\psi_{h_{1}}$ is uniquely determined. Furthermore any phase $\psi_{\mathbf{h}_{2}}$ which is lincarly semi-independent of $q_{\mathbf{h}_{1}}$ then has two possible values differing from each other by $\pi$. Either one of these two values for a particular such phase $\psi_{h_{2}}$ may be chosen. Once this is done the value of any phase $\gamma_{\mathbf{h}}$, of necessity linearly semi-dependent on $q_{\mathbf{h}_{\underline{2}}}$, is uniquely determined.

## 4•10. Type $3 P_{3}{ }^{0}$

Theorem $4 \cdot 10 \cdot 1$. A single phase $\psi_{h}$ is a seminvariant if, and only if, $h+k=0$.

Theorem $4 \cdot 10 \cdot 2$. Let $h_{1}+k_{1} \neq 0$, so that $\varphi_{\mathbf{h}_{1}}$ is linearly semi-independent. Then the value of $q_{h_{1}}$ may be
specified arbitrarily. Once this is done, the value of any phase $\varphi_{h}$ which is linearly semi-dependent on $\varphi_{\mathbf{h}_{1}}$ is uniquely determined. Any phase $\varphi_{h}$, of necessity rationally semi-dependent on $\psi_{h_{1}}$, is also linearly semi-dependent on $\psi_{h_{1}}$, provided that $\psi_{h_{1}}$ is semiprimitive, i.e. provided that $h_{1}+k_{1}= \pm 1$.

## 4•11. Type $3 P_{3}{ }^{2 *}$

Theorem $4 \cdot 11 \cdot 1$. A single phase $\psi_{h}$ is a seminvariant if, and only if, $h+k$ is even.

Theoren $4 \cdot 11 \cdot 2$. Let $h_{1}+k_{1}$ be odd so that $\varphi_{h_{1}}$ is linearly semi-independent. Then $\varphi_{\mathbf{h}_{1}}$ has just two possible values and these differ from each other by $\pi$. Either one of these two values may be chosen. Once this is done, the value of any phase, of necessity linearly semi-dependent on $\mathscr{f}_{\mathbf{h}_{1}}$ is uniquely determined.

## 4•12. Type $3 P_{3} 4$

Theorem 4•12•1. A single phase $\psi_{h}$ is a seminvariant if, and only if, $h-k \equiv 2 l(\bmod 4)$.

Theorem $4 \cdot 12 \cdot 2$. Let the phase $\psi_{h_{1}}$ be linearly semiindependent. Depending upon whether $h_{1}-k_{1}+2 l_{1}$, is odd or even, there are four or two possible values for $\psi_{h_{1}}$ (differing by $\pi / 2$ or $\pi$, respectively).

In the first case any of the four possible values for $\psi_{h_{1}}$ may be chosen. Once this is done, the value of any phase $\psi_{h}$, of necessity linearly semi-dependent on $\varphi_{\mathbf{h}_{1}}$, is uniquely determined.

In the second case either of the two possible values for $\mathscr{F}_{h_{1}}$ may be chosen. Once this is done, the value of any phase $\psi_{h}$ which is linearly semi-dependent on $\vartheta_{h_{1}}$ is uniquely determined. Furthermore any phase $\boldsymbol{q}_{h_{2}}$ which is linearly semi-independent of $\mathscr{F}_{\mathbf{h}_{1}}$, then has two possible values differing from cach other by $\pi$. Either one of these two values, for a particular such phase $\psi_{h_{2}}$, may be chosen. Once this is done the value of any phase $q_{h}$, of necessity linearly semi-dependent on $q_{h_{2}}$, is uniquely determined.

## 4•13. Type $3 P_{4}{ }^{\prime}$ )

Theorem $4 \cdot 13 \cdot 1$. A single phase $\psi_{h}$ is a seminvariant if, and only if, $h+k=l$.

Theorem $4 \cdot 13 \cdot 2$. Let $h_{1}+k_{1}-l_{1} \neq 0$, so that $\varphi_{\mathbf{h}_{1}}$ is linearly semi-independent. Then the value of $\varphi_{h_{1}}$ may be specified arbitrarily. Once this is done, the value of any phase $q_{h}$ which is linearly semi-dependent on $q_{h_{1}}$ is uniquely determined. Any phase $q_{\mathbf{h}}$, of necessity rationally semi-dependent on $\mathscr{q}_{h_{1}}$, is also linearly semidependent on $q_{\mathbf{h}_{1}}$, provided that $\gamma_{\mathbf{h}_{1}}$ is semi-primitive, i.e. provided that $h_{1}+k_{1}-l_{1}= \pm 1$.

## 4•14. Type 4PIII

Theorem 4•14•]. Every phase is a seminvariant.

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## 5. Concluding remarks

This paper concludes the study of the seminvariants for the non-centrosymmetric space groups which was initiated in a previous paper (Hauptman \& Karle, 1956). The theory of the seminvariants provides a basis for specifying an origin and the enantiomorph or reference frame when required. Furthermore it demonstrates the existence of relationships between the measured intensities and the values of phases. It will be the purpose of future publications to elucidate the exact nature of these relationships and by these means to continue the unified program for phase determina-
tion in the non-centrosymmetric space groups which has already been completed for the centrosymmetric. ones (Karle \& Hauptman, 1961 ff.).

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# Neutron Diffraction Investigation of Solid Solutions $\mathbf{A l T h}_{2} \mathbf{D}_{\boldsymbol{n}}$ 

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Solid solutions of composition $\mathrm{AlTh}_{2} \mathrm{D}_{n}$, with $n=0,2,3,4$, have been studied by means of neutron diffraction. For $n=4$ the deuterium atoms completely fill a set of equivalent Th-tetrahedra, quite similar to the arrangement in thorium hydride. For the other compositions these sites are partly occupied. No evidence for ordering has been found, even at a temperature of $8 \mathbf{2} \mathbf{~ " K}$.

The intermetallic compound AITh ${ }_{2}$ easily absorbs hydrogen. Apart from a two-phase region at room temperature between the compositions $\mathrm{AlTh}_{2} \mathrm{H}_{0}$ and $\mathrm{AlTh}_{2} \mathrm{H} \sim_{1.5}$, the hydrogen is dissolved homogeneously until the ultimate composition $\mathrm{AlTh}_{2} \mathrm{H}_{4}$ is reached (van Vucht, 1960). X-ray investigation shows that the tetragonal symmetry of AlTh. is conserved in the solid solutions. When the lattice parameters are plotted against $n$, the number of hydrogen atoms per $\mathrm{AlTh}_{2}, a$ is found to increase up to $n=2$. There it shows a sharp break, followed by a decrease until saturation. On the other hand $c$ increases monotonically.

As part of a larger program, a neutron-diffraction investigation was undertaken with the object of establishing the hydrogen positions. Only microcrystalline samples were available so that to avoid a large background of incoherent scattering the deuterides rather than hydrides were used. The relevant neutron scattering lengths (Shull \& Wollan, 1956) are, in $10 \cdot 12 \mathrm{~cm} ., b_{\mathrm{Al}}=0 \cdot 35, b_{\mathrm{Th}}=1 \cdot\left(0 \mathrm{I}\right.$ and $b_{\mathrm{I}}=0.65$.

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## Experimental procedure

The deuterides were prepared in exactly the same way as the hydrides (van Vucht, 1960). For the roomtemperature neutron-diffraction measurements 10 mm . dia. cylindrical thin-walled aluminium sample holders were used. By means of a glass tube and a section of fernico tube these were connected to the apparatus in which the deuteride was prepared. Using a tilting arrangement the finished product could be transferred to the sample holder under vacuum after which the glass connecting tube was sealed off. The sample holder was then placed on the diffractometer described by Goedkoop (1957) and the diffraction pattern recorded with $1 \cdot(126$ A neutrons. Resolution was mainly determined by Soller slits 0.25 mm . wide and 200 mm . long placed in front of the counter.

For measurements at low temperature a singlejacketed vacuum cryostat as shown in Fig. I was placed on the goniometer. Liquid air or liquid nitrogen was placed in the inner cylinder, to the bottom of which the sample holder was fixed. The glass-sealed sample holders were unsuited for this arrangement and so a shorter one closed by means of a screw-plug


[^0]:    * For space group $P\left(14_{1} 22\right)$, the signs of all seminvariants are uniquely determined. In this case, therefore, the specification of the sign of a seminvariant is not a requirement for theorems $4 \cdot 11 \cdot 1$ and $4 \cdot 11 \cdot 2$ to be valid.

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